

# Heavy Quark Flows as Better Probes of QGP Properties

Zi-Wei Lin  
East Carolina University

Based on Hanlin Li, ZWL, Fuqiang Wang, in preparation;  
and  
ZWL's talk at Strange Quark Matter 2017

# Outline

- Introduction of a multi-phase transport (AMPT) model
- Anisotropic parton escape: *a brief review*
- Anisotropic parton escape: *flavour dependence*
- Summary

# For comprehensive simulations of high energy heavy ion collisions

## We need:

Initial particle/energy production



Pre-equilibrium interactions:  
*equilibration, thermalization, initial flow*



Space-time evolution of QGP



Hadronization  
/QCD phase transition



Hadronic interactions

## Choices:

*Soft+hard model (such as HIJING),  
CGC, pQCD, ...*



*Parton cascade (ZPC, MPC, BAMPS),  
NJL, CGC, AdS/CFT, ...*



*Parton cascade (ZPC, MPC, BAMPS), NJL,  
(ideal, viscous, anisotropic) hydrodynamics, ...*



*Quark coalescence/parton recombination,  
string fragmentation, Cooper-Frye, statistical  
hadronization, independent fragmentation,  
rate equations, ...*

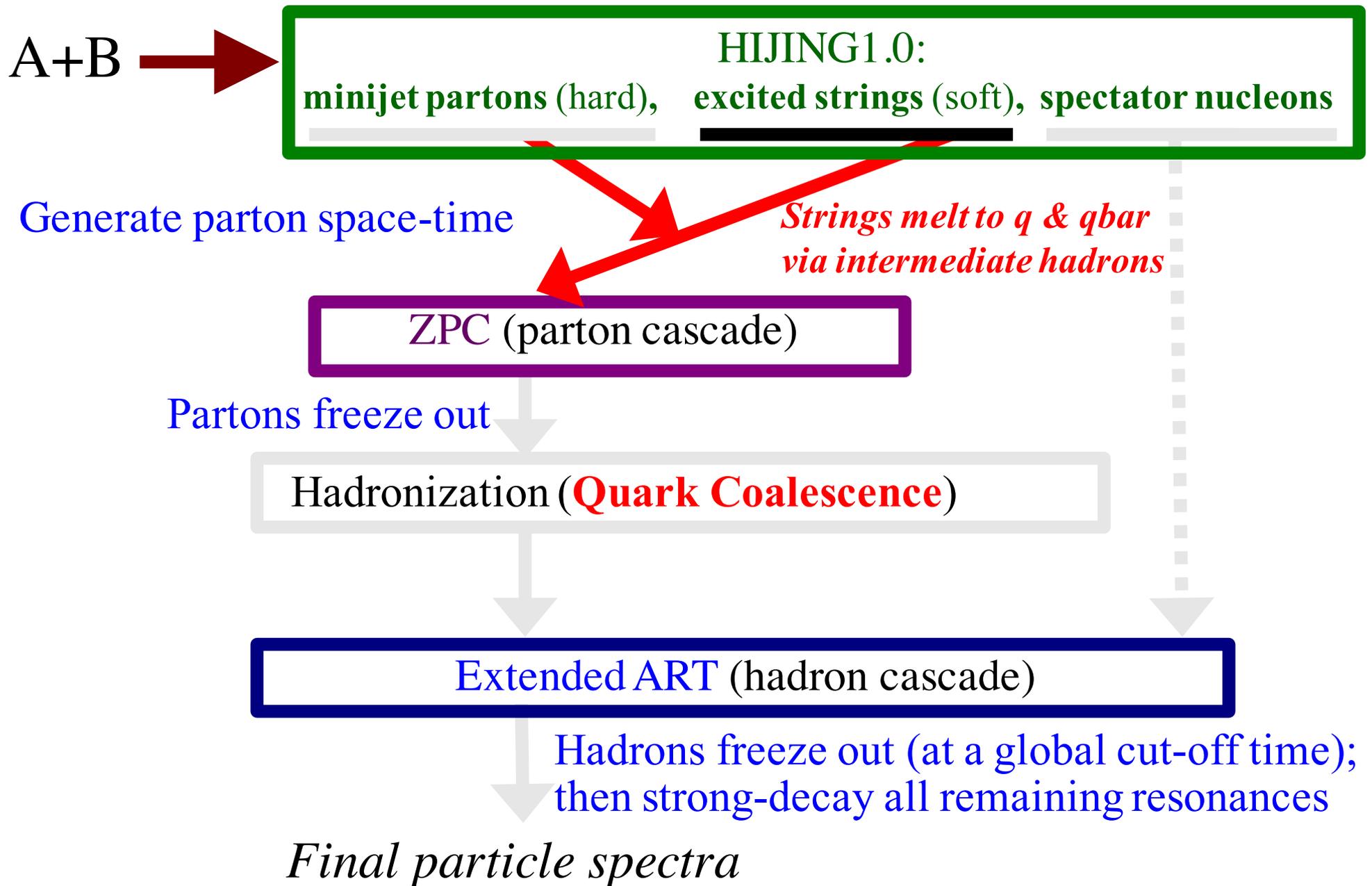


*Hadron cascade (ART, RQMD, UrQMD, ...),  
thermal model (w/ freezeout temperatures), ...*

The AMPT model currently includes the **green** components for each phase.

ZWL et al. PRC72 (2005).

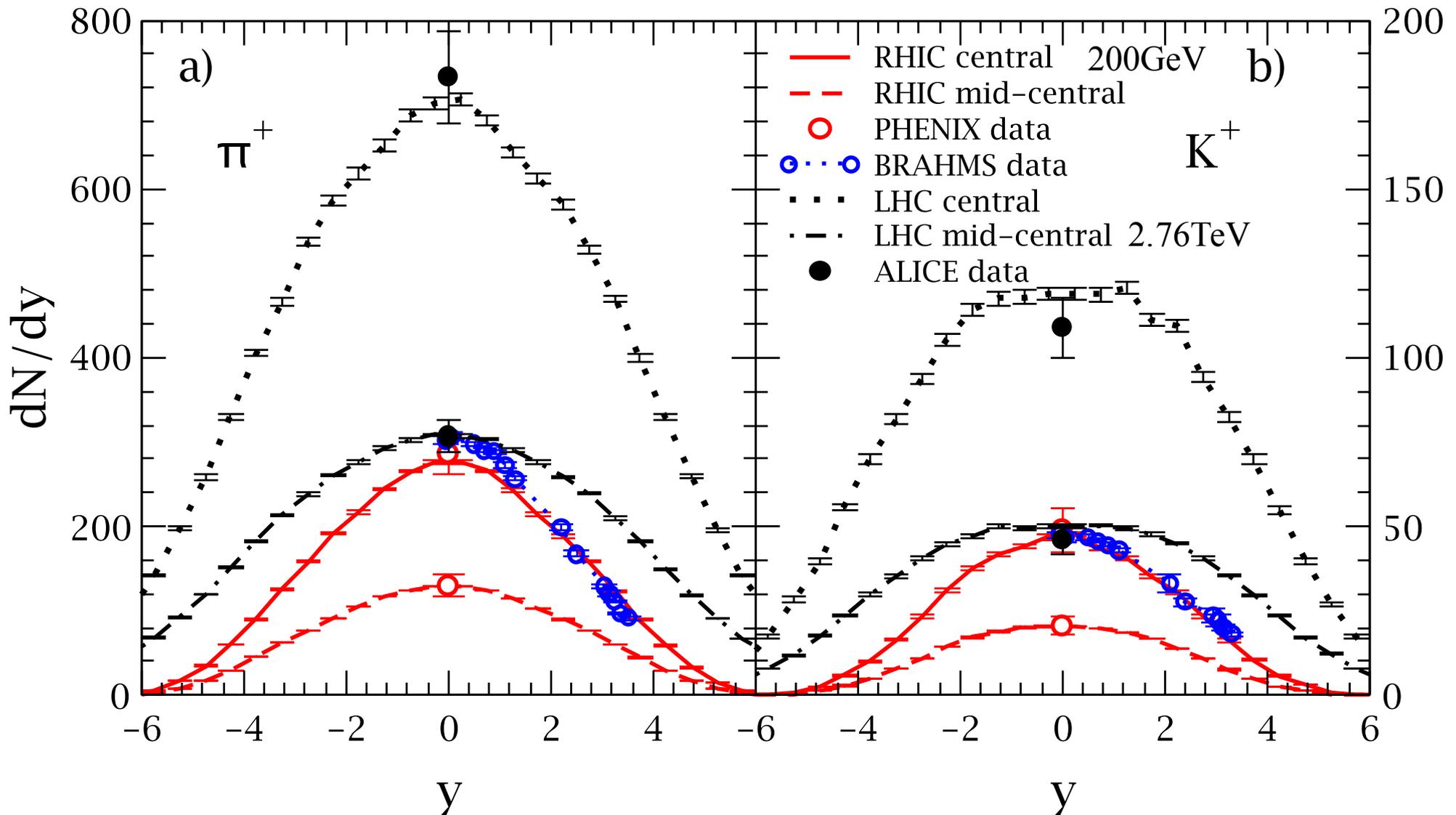
# Structure of AMPT v2.xx (String Melting version)



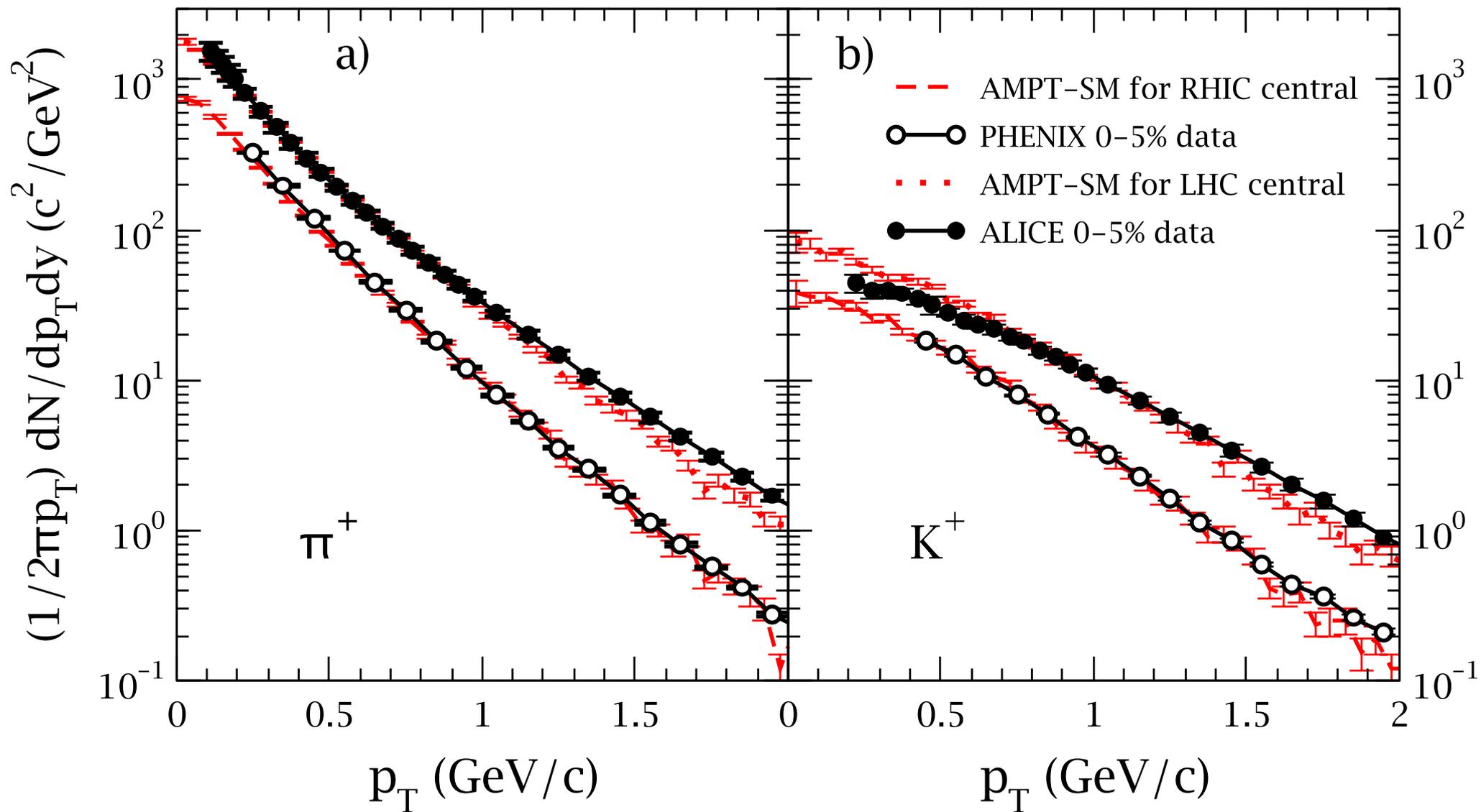
Now the String Melting AMPT can reasonably describe the bulk matter at high energies at RHIC and LHC.

$dN/dy$  of  $\pi$  & K:

ZWL, PRC 90 (2014)

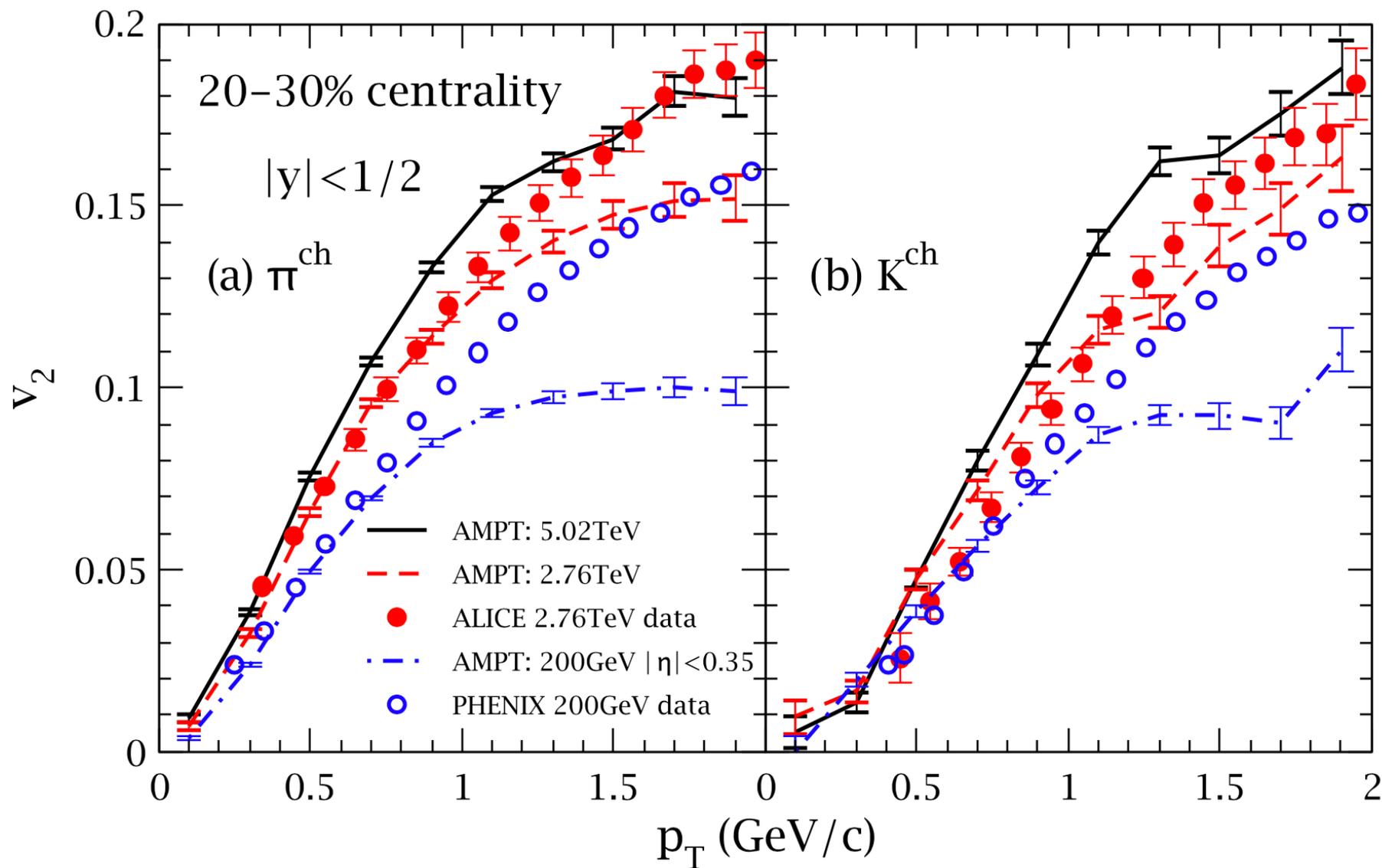


$p_T$ -spectra of  $\pi$  & K (in central collisions): ZWL, PRC 90 (2014)

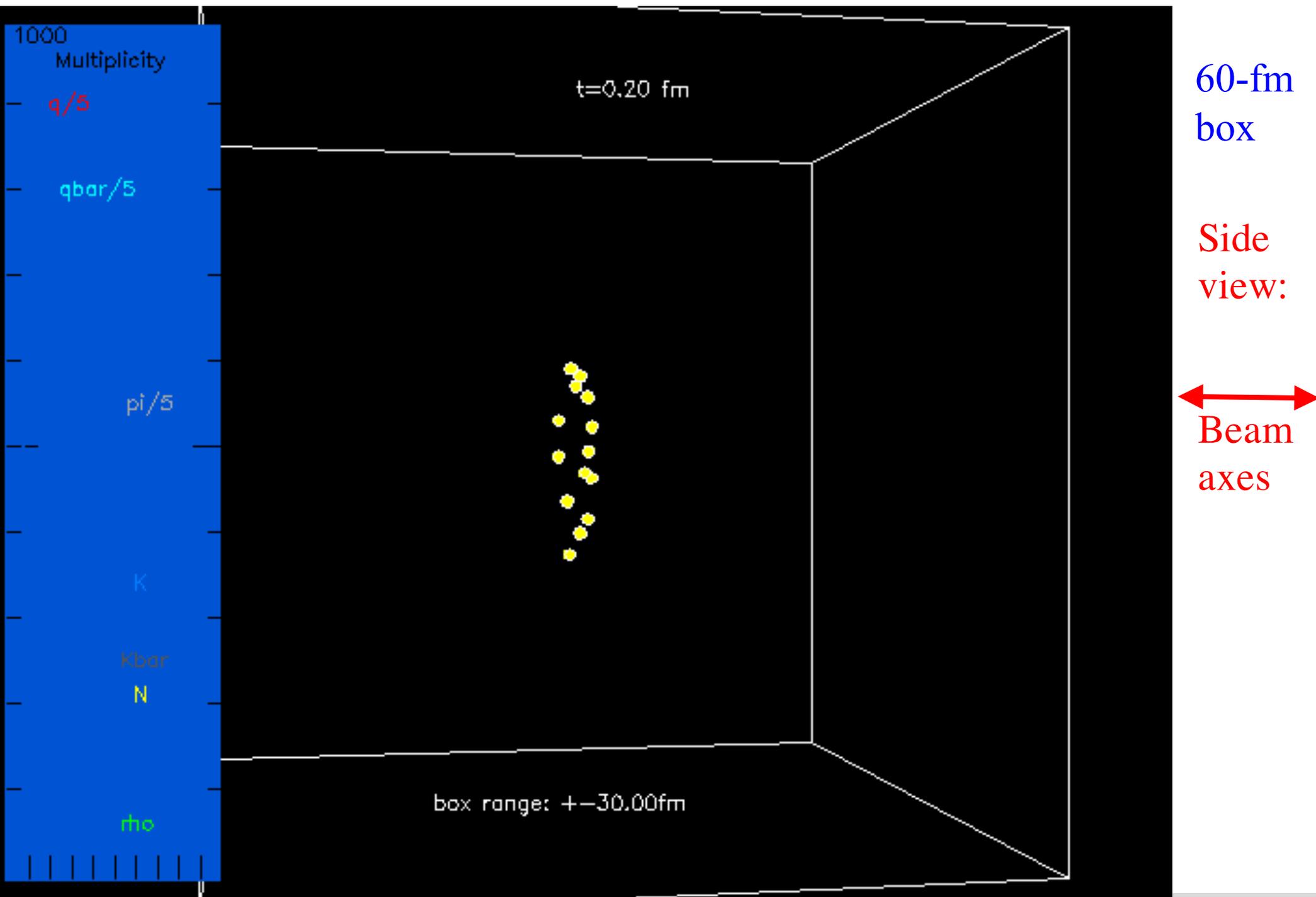


$v_2$  of  $\pi$  & K (in mid-central collisions):

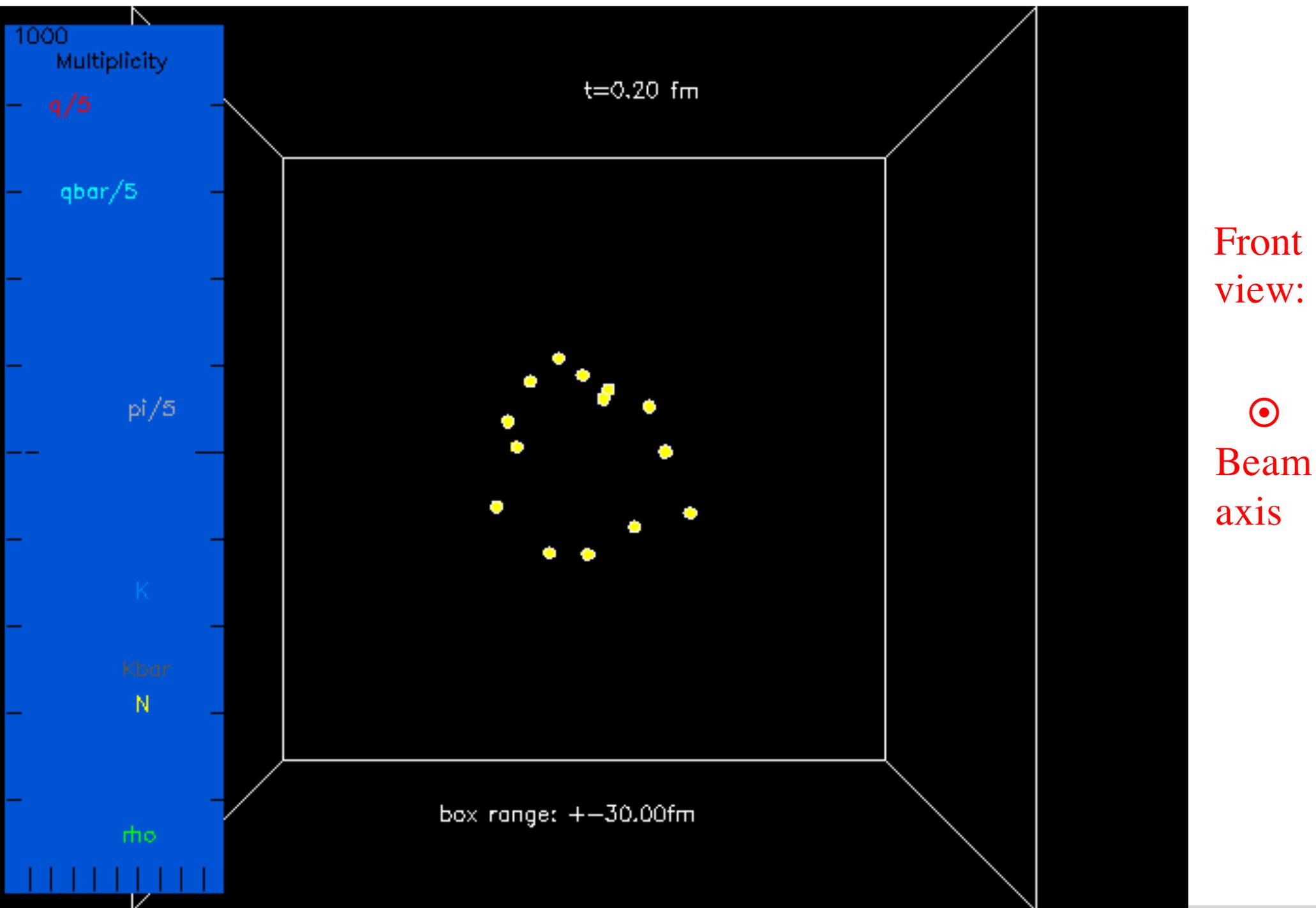
Guo-Liang Ma & ZWL,  
PRC 93 (2016)



# One central Au+Au event at 200A GeV from String Melting AMPT



# The same Au+Au event from a different viewpoint



Current and previous public versions of AMPT are available at

<http://myweb.ecu.edu/linz/ampt/>



## AMPT source codes

(updated October 28, 2016):

A Multi-Phase Transport (AMPT) model is a Monte Carlo transport model for nuclear collisions at relativistic energies.

Each of the following versions contains:

**the source codes, an example input file, a Makefile, a readme, a required subdirectory for storing output files, and a script to run the code.**

1. [ampt-v1.11-v2.11.tgz](#) (11/2004)
2. [ampt-v1.21-v2.21.tgz](#) (10/2008)
3. [ampt-v1.25t3-v2.25t3.tgz](#) (8/2009)
4. [ampt-v1.25t7-v2.25t7.zip](#) (9/2011)
5. [ampt-v1.25t7d-v2.25t7d.zip](#) (4/2012)
6. [ampt-v1.26t1-v2.26t1.zip](#) (9/2012)
7. [ampt-v1.26t4-v2.26t4.zip](#) (8/2014)
8. [ampt-v1.26t5-v2.26t5.zip](#) (4/2015)
9. [ampt-v1.26t7-v2.26t7.zip](#) (10/2016)

**This readme file lists the main changes up to version v1.26t7-v2.26t7 ("t" means a version under test):**

AMPT Users' Guide

\*\*\*\*\*

5/2016 test version v1.26t7/v2.26t7:

\* (w/ GL Ma) Changed the following line in art1f.f to avoid the bug

# The escape mechanism: *a brief review*

Liang He, Terrence Edmonds, ZWL, Feng Liu, Denes Molnar, Fuqiang Wang: PLB 753 (2016):

Anisotropic parton escape is the dominant source of azimuthal anisotropy in transport models.

ZWL et al. NPA 956 (2016) for Quark Matter 2015:

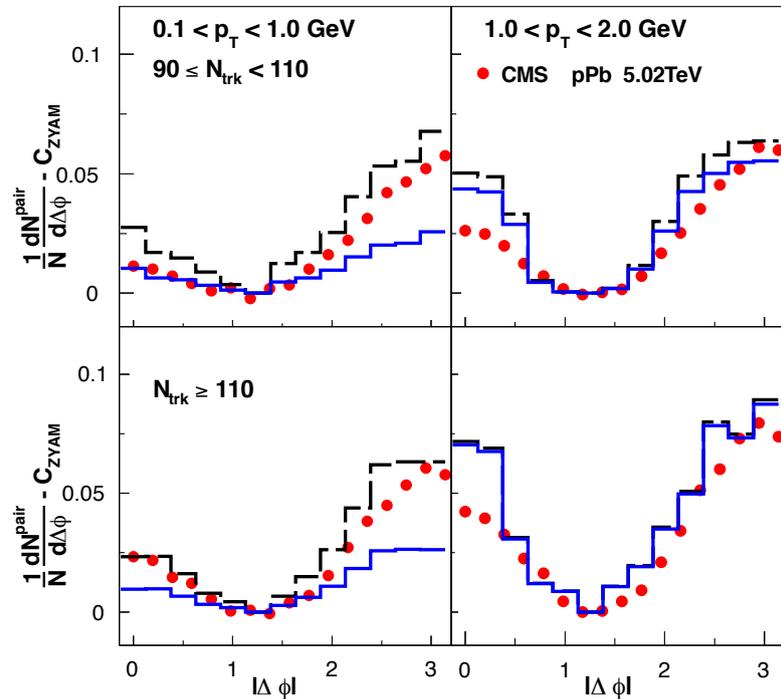
Elliptic anisotropy  $v_2$  may be dominated by particle escape instead of hydrodynamic flow.

**It has been generally believed that:**

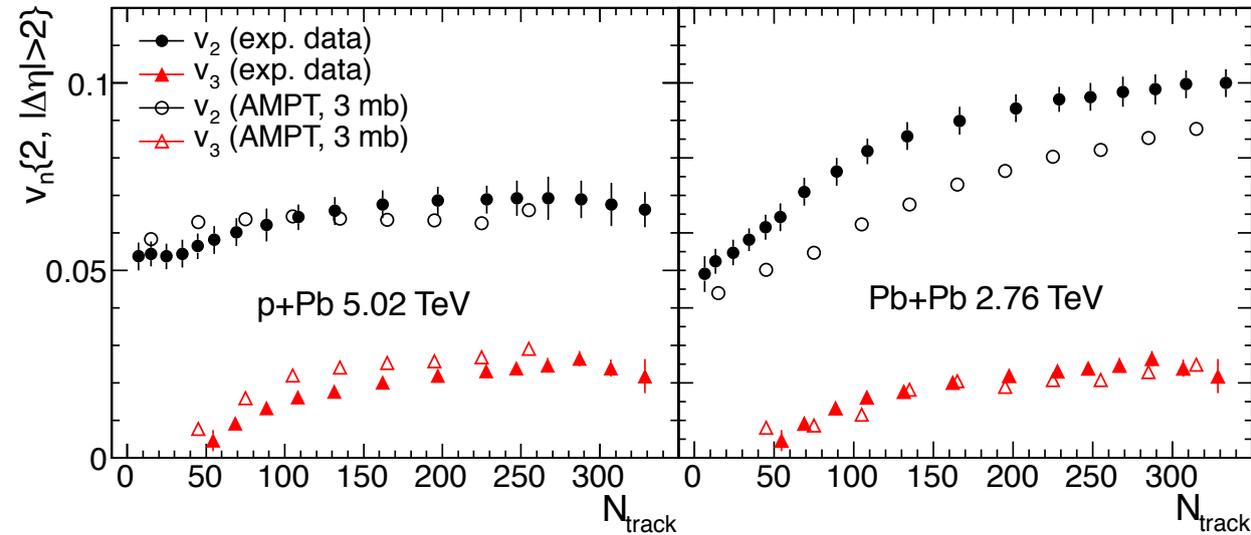
- **Transport models** at large-enough cross section will approach **hydrodynamics**.
- Early hydro-type collective flow in sQGP converts initial spatial anisotropy into final momentum-space  $v_n$
- For low- $P_T$  particles in high-energy heavy ion collisions, since both **hydrodynamics** and **transport models** can describe  $v_n$  data, the mechanism of  $v_n$  development in **transport models** (*via particle interactions*) is in principle the same as in **hydrodynamics** (*via pressure gradients*).

# The escape mechanism: *a brief review*

Small systems: again, both **hydrodynamics** and **transport** can describe flow.



Bozek and Broniowski, PLB 718 (2013)  
using e-by-e viscous hydrodynamics.



Bzdak and Ma, PRL 113 (2014)  
using AMPT (String Melting version).

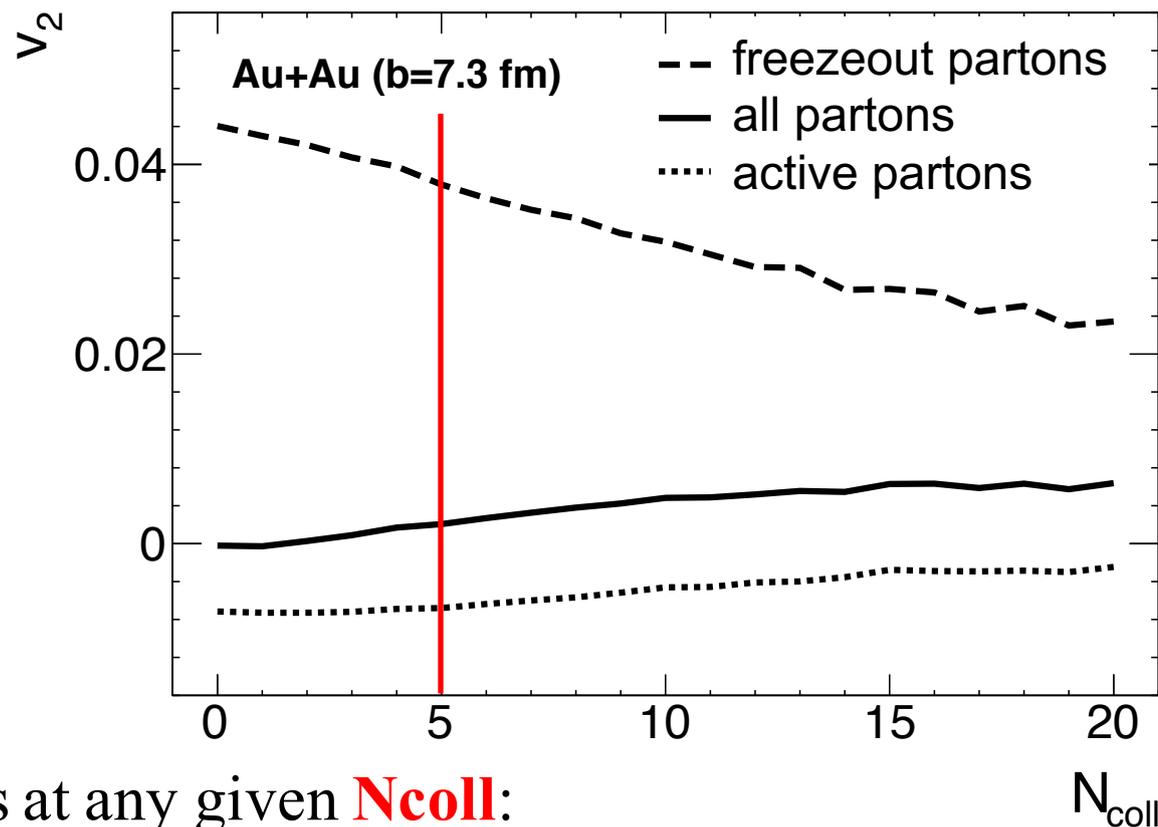
## Questions for small systems such as p+Pb or d+Au:

- Mean free path may be comparable to the system size;  
is **hydrodynamics** still applicable to such small systems?
- Transport and hydrodynamics should be different for small systems,  
could they also be different for large systems?

## The escape mechanism: *a brief review*

We have followed the complete parton collision history and study the generation of parton  $v_2$  in AMPT. He et al. PLB753 (2016)

**Ncoll**: *number of collisions suffered by a parton.*



3 parton populations at any given **Ncoll**:

$N_{\text{coll}}$

- freezeout partons:** *freeze out/hadronize after exactly  $N_{\text{coll}}$  collisions;*
- active partons:** *will collide further, freeze out after  $>N_{\text{coll}}$  collisions;*
- all partons:** *sum of the above two populations (i.e. all partons that have survived  $N_{\text{coll}}$  collisions).*

# The escape mechanism: *a brief review*

**At  $N_{\text{coll}}=0$ :**

all partons:  $v_2=0$

they contain 2 parts:

freezeout:  $v_2 \approx 4.5\%$ ,

active:  $v_2 < 0$ .

by symmetry (*since they include all initial partons*);

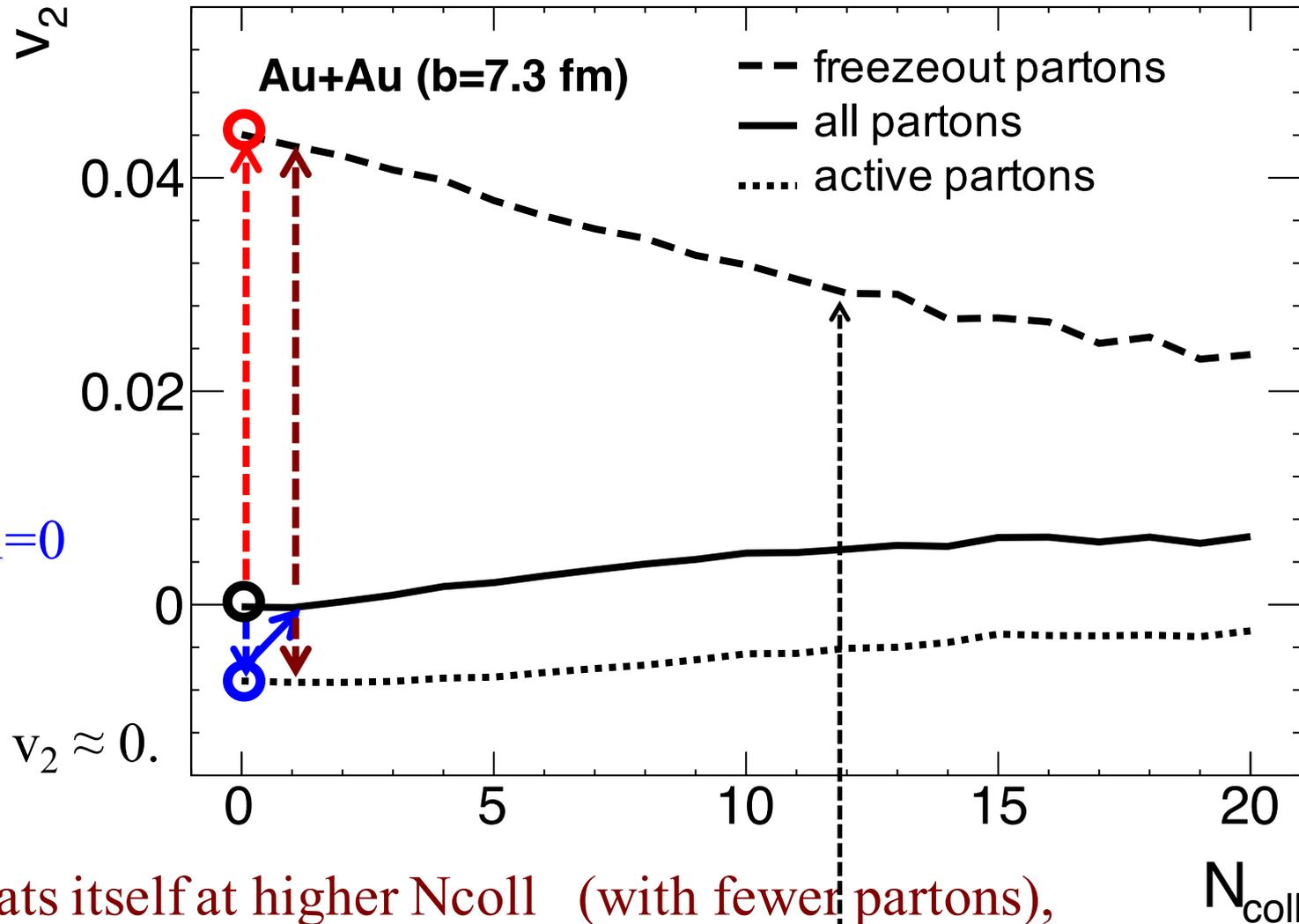
**At  $N_{\text{coll}}=1$ :**

active partons at  $N_{\text{coll}}=0$

collide once each

& become

all partons at  $N_{\text{coll}}=1$ :  $v_2 \approx 0$ .



This process repeats itself at higher  $N_{\text{coll}}$  (with fewer partons),  
eventually all partons freezeout/hadronize.

$\langle v_2 \rangle$  = weighed average of the freezeout partons'  $v_2$  at different  $N_{\text{coll}}$ .

# The escape mechanism: *a brief review*

**At  $N_{\text{coll}}=0$ :**

freezeout partons:  $v_2 \approx 4.5\%$ ,

this is **purely** due to the

**anisotropic escape probability** (*escape mechanism*):

interaction-induced response to geometrical shape,  
no contribution from collective flow.

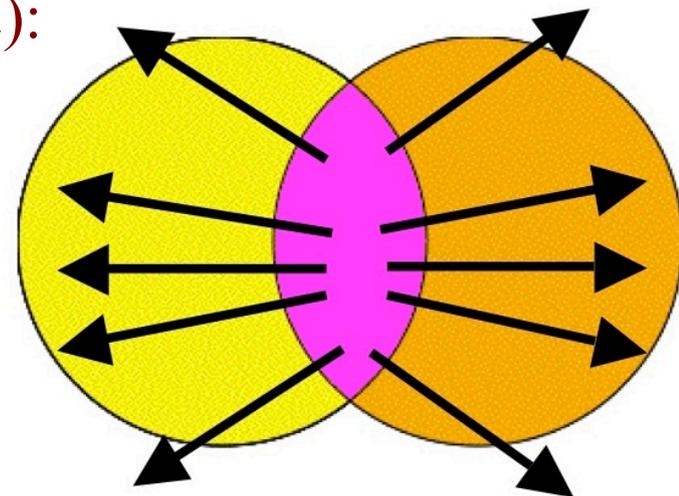
**At  $N_{\text{coll}} \geq 1$ :**

freezeout partons:  $v_2 > 0$

due to

**anisotropic escape probability**  
& **(anisotropic) collective flow**.

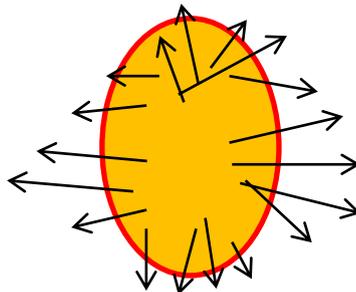
In event-averaged picture  
of elliptic flow:



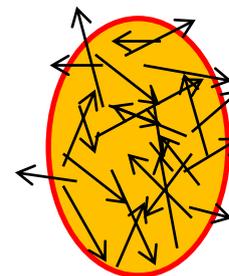
How to separate the two contributions?

We design a **Random- $\phi$  Test** (destroy collective flow but keep the anisotropic shape):

**Normal:**



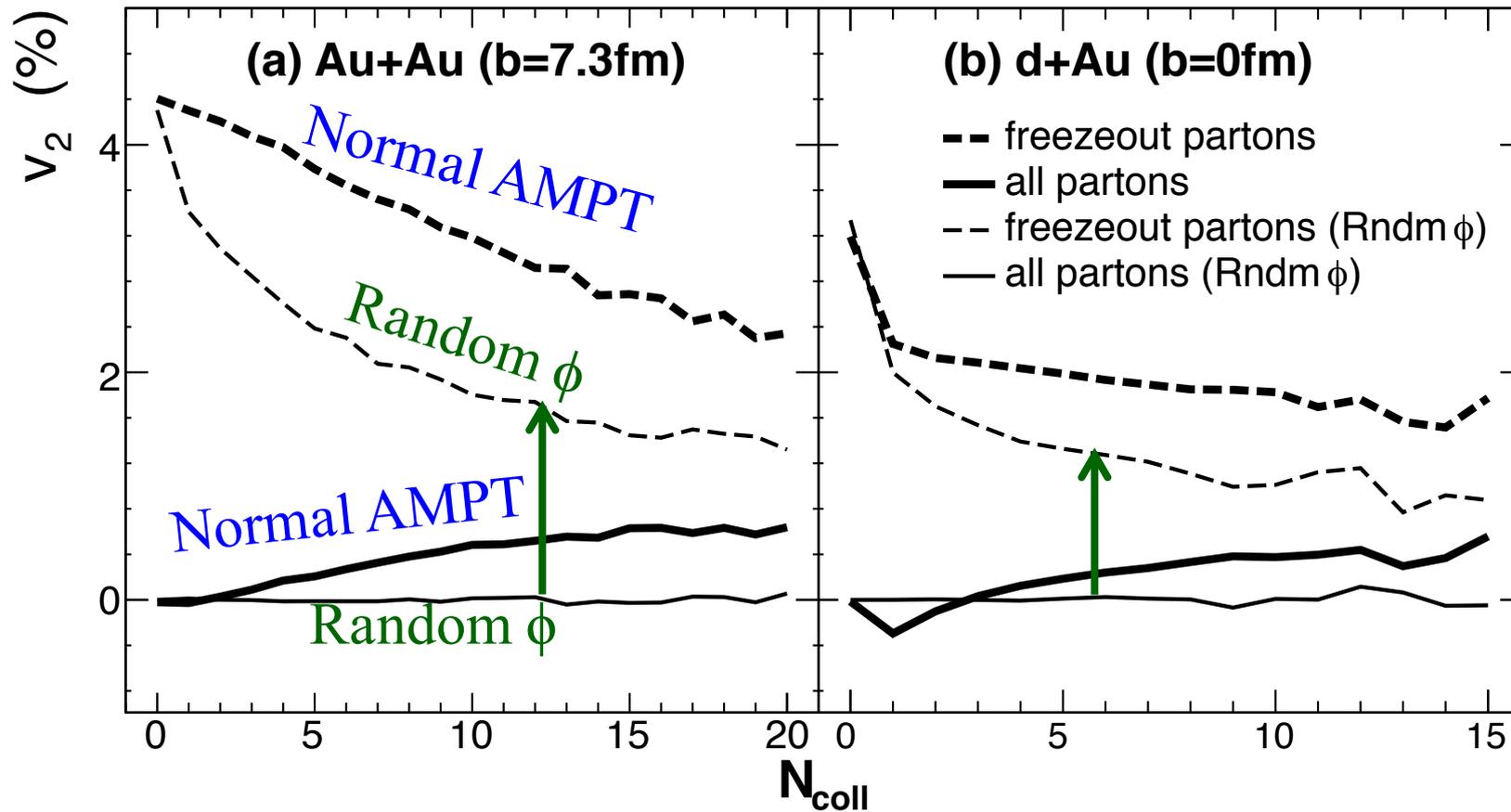
**Random:**



# The escape mechanism: *a brief review*

$v_2$  from the **Random- $\phi$  Test**: purely from the escape mechanism

He et al. PLB753 (2016)



$\langle v_2 \rangle_{\text{normal}}$      $\langle v_2 \rangle_{\text{random-}\phi}$     **Ratio of random/normal**     $\langle N_{\text{coll}} \rangle$   
 ~ fraction from pure escape

Au+Au	3.9%	2.7%	<b>69%</b>	4.6 ( <i>modest</i> )
d+Au	2.7%	2.5%	<b>93%</b>	1.2 ( <i>low</i> )

# The escape mechanism: *a brief review*

Elliptic anisotropy  $v_2$  may be dominated by particle escape instead of hydrodynamic flow.

## Implications:

- The escape mechanism helps to explain similar anisotropic flows observed in small and large systems:  
since both are dominated by the same mechanism (*anisotropic escape probability*)
- The driving force for  $v_2$  at low & high  $P_T$  is qualitatively the same  
since both are dominated by *anisotropic probability of interactions* before escape  
(scatterings/kicks for low  $P_T$  & energy loss for high  $P_T$ )
  - At **low-to-modest** opacity or  $\langle N_{\text{coll}} \rangle$ :  
transport and hydrodynamics are different;  
the escape mechanism dominates  $v_n$ .
  - At **very high** opacity or  $\langle N_{\text{coll}} \rangle$ :  
transport and hydrodynamics are similar;  
hydro-type collective flow dominates  $v_n$ .

# The escape mechanism: *flavour dependence*

These previous results are for **all quarks @200 GeV**.

**Q:**  
**Does the escape mechanism work differently for different flavours?**  
**or**  
**Does collective flow work differently for different flavours?**

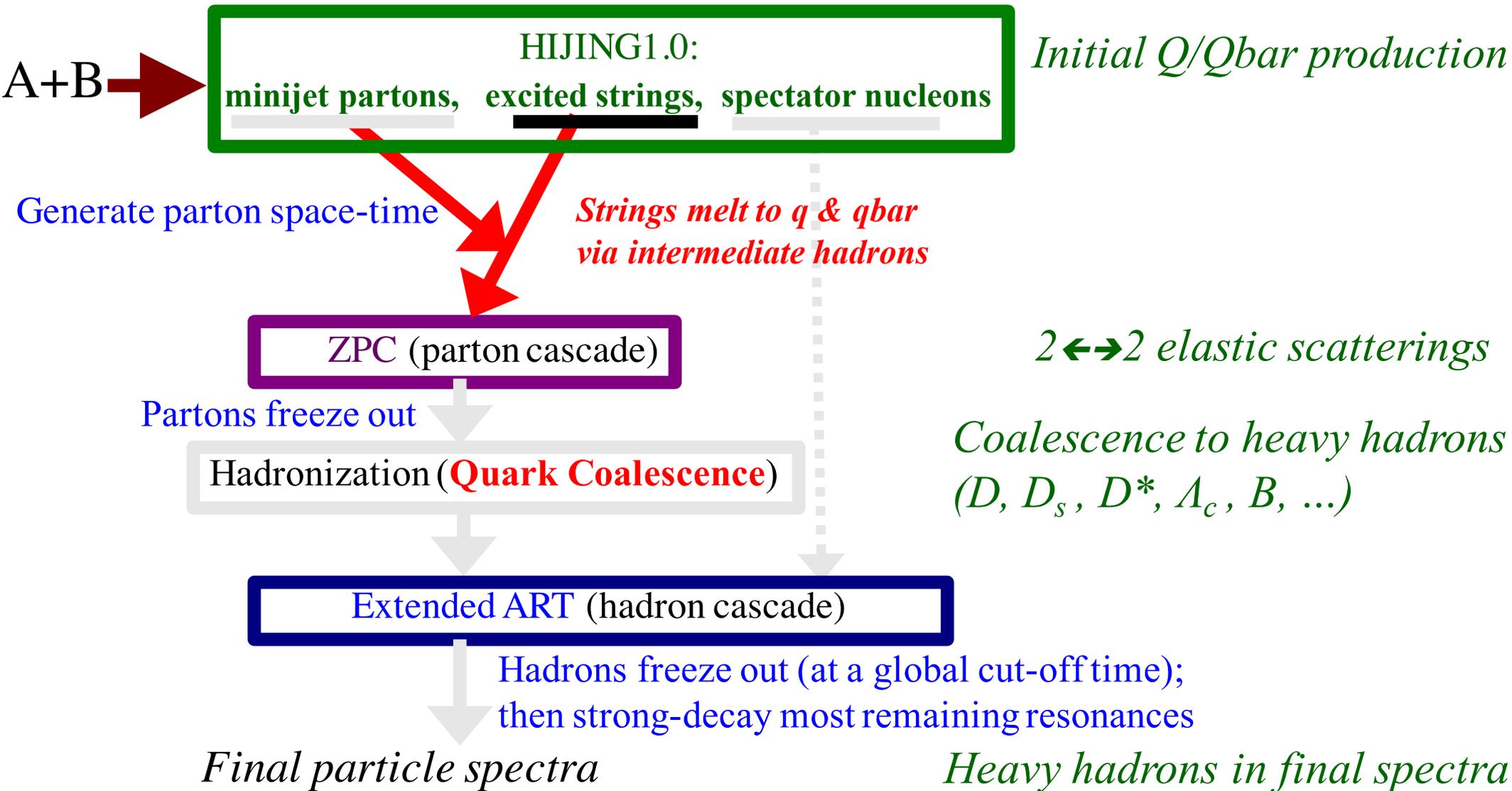
We now use string melting AMPT to analyze  
**light (u/d), strange, charm quarks**  
in p+Pb@5TeV, Au+Au@200GeV, Pb+Pb@2.76TeV.  
Hanlin Li, ZWL, Fuqiang Wang, in preparation.

3mb parton cross section is used      ZWL, PRC 90 (2014)  
since it reproduces pi/K/p  $v_2(P_T)$ .      G.L. Ma & ZWL, PRC 93 (2016)

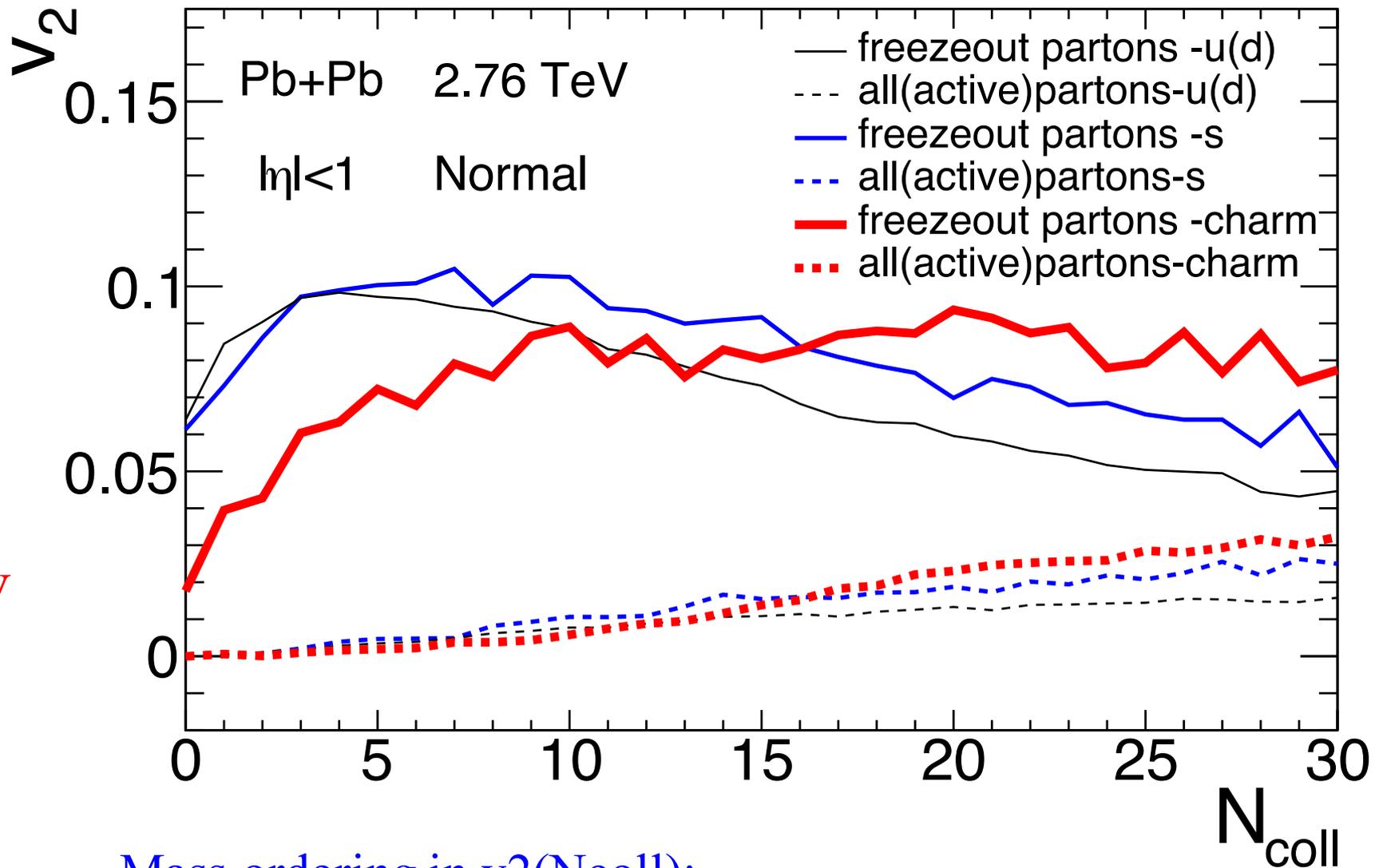
*Caveat: here we use the same cross section for all flavours.*

# Where are heavy flavours (Q) in the current AMPT model?

## Structure of AMPT v2.xx (String Melting version)



# The escape mechanism: *flavour dependence*

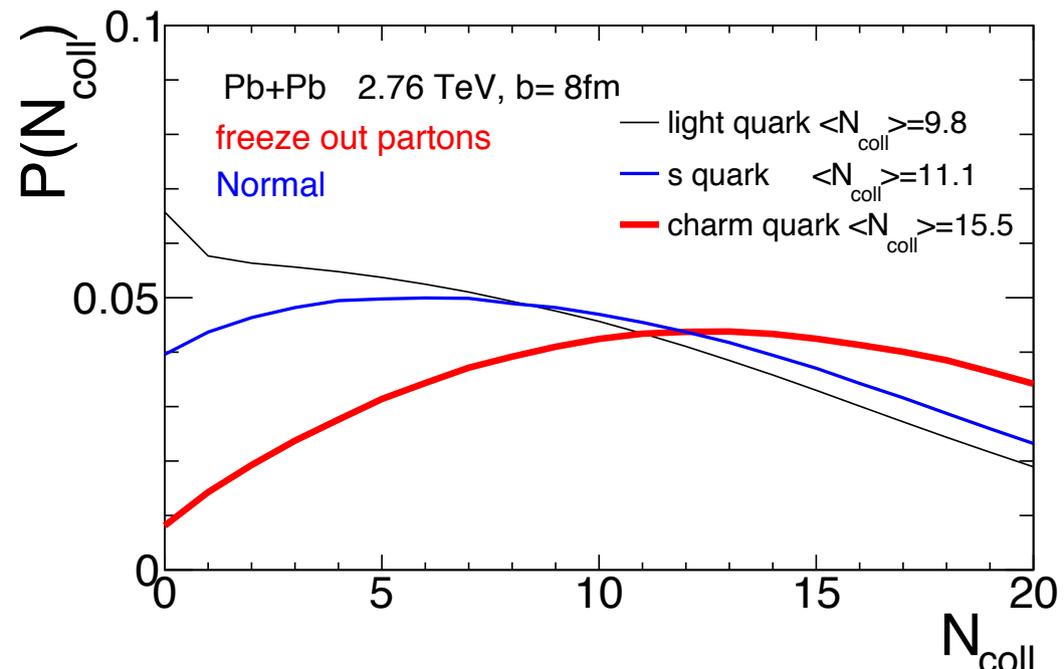
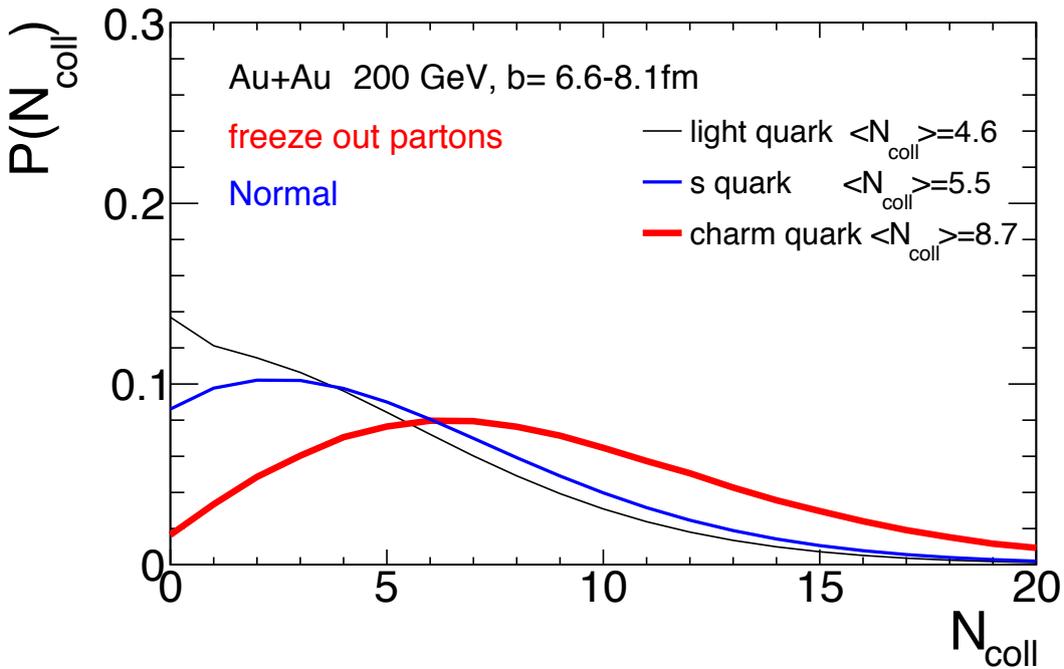


Mass ordering in  $v_2(N_{\text{coll}})$ :

$v_2^c < v_2^s < v_2^{ud}$  at small  $N_{\text{coll}}$ ;

reversed:  $v_2^c > v_2^s > v_2^{ud}$  at large  $N_{\text{coll}}$ .

# The escape mechanism: *flavour dependence*

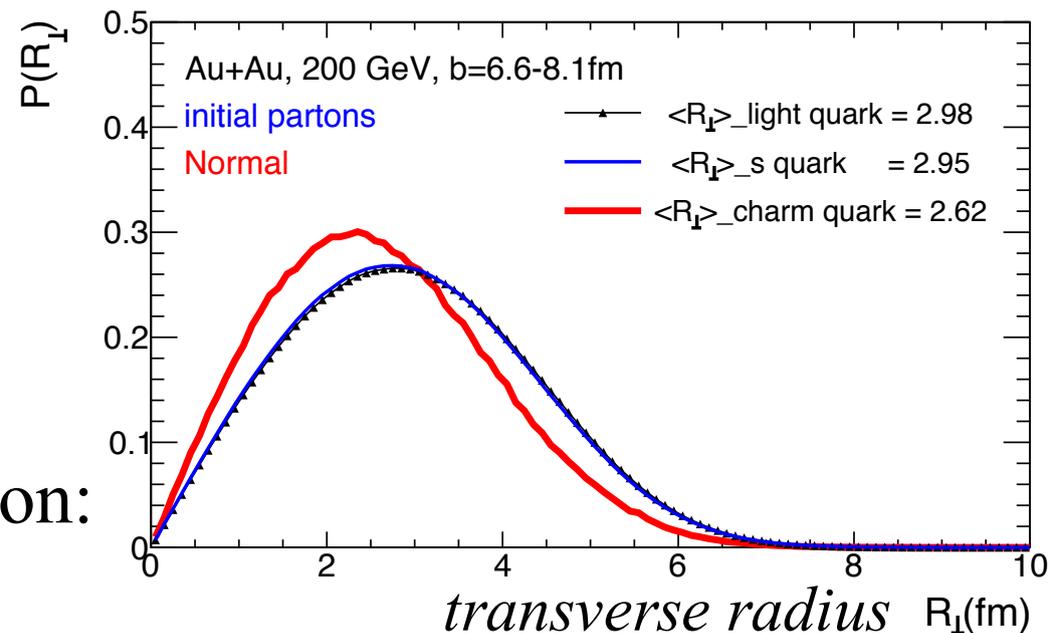


Mass ordering in the  $N_{\text{coll}}$  distribution

for all 3 systems:

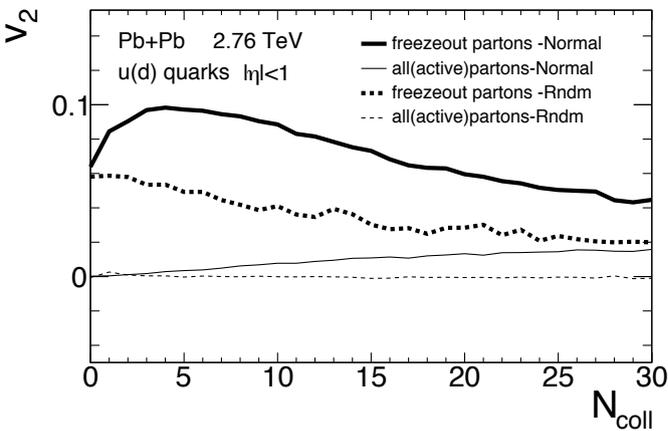
$$\langle N_{\text{coll}} \rangle_{\mathbf{c}} > \langle N_{\text{coll}} \rangle_{\mathbf{s}} > \langle N_{\text{coll}} \rangle_{\mathbf{ud}}$$

This is related to the initial  
(velocity &) spatial distribution:

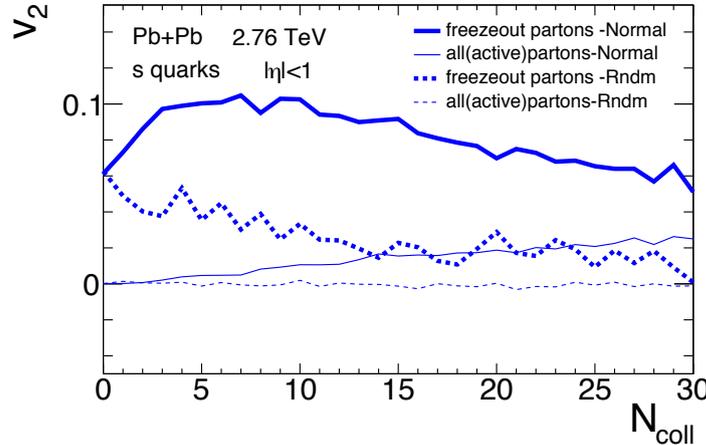


# The escape mechanism: *flavour dependence*

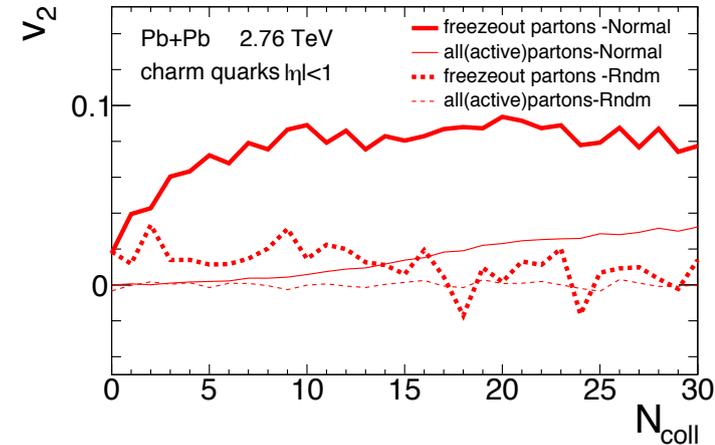
u/d



s



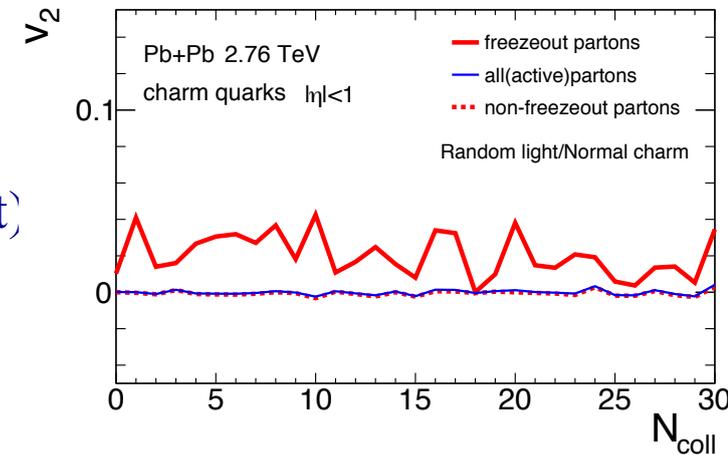
c



Pb+Pb  
2.76TeV  
8fm

Random- $\phi$  test (*randomized quarks of all flavours*):  
shows greater reduction of  $v_2$  for heavier quarks.

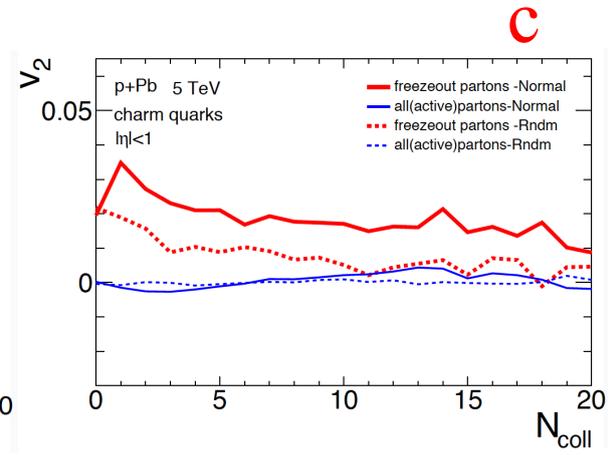
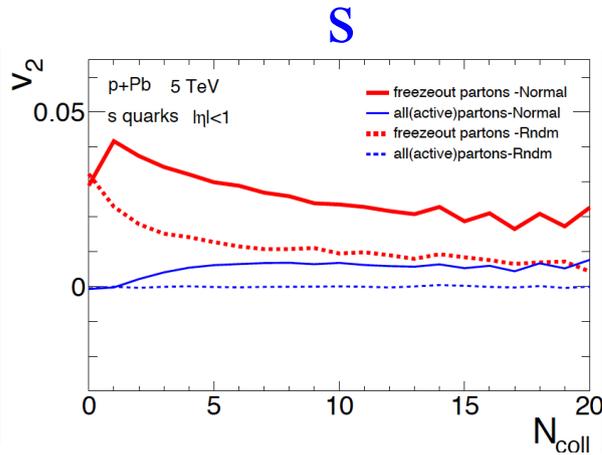
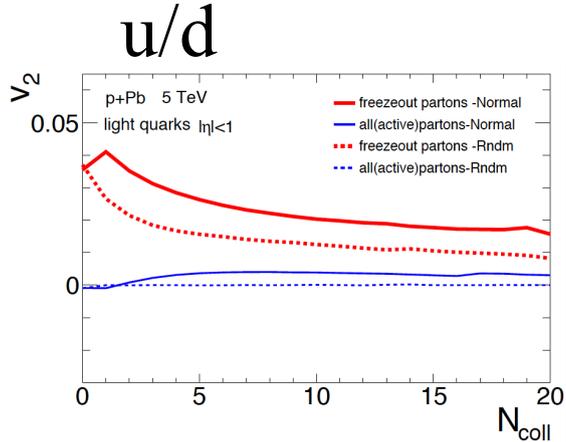
We also tested: *random- $\phi$  for u/d/s quarks only*  
(*normal charm: charm quarks keep their collective flow*):  
large reduction of charm  $v_2$  (like all-flavour random- $\phi$  test)  
→  
light quark collective flow is essential for charm  $v_2$



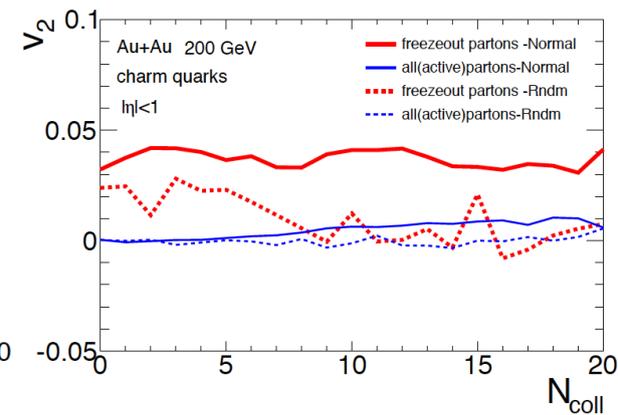
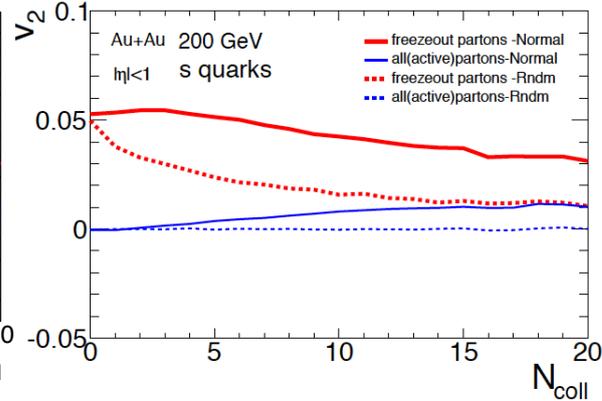
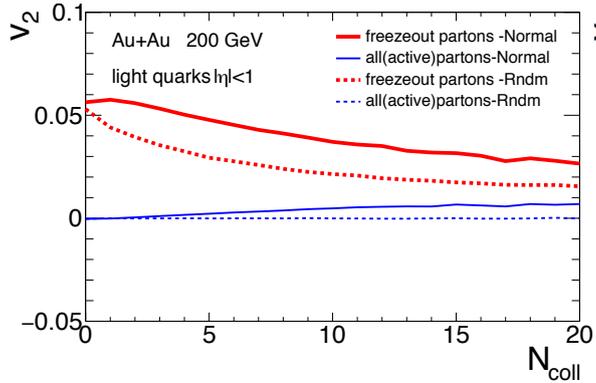
# The escape mechanism: *flavour dependence*

Analysis is done for 3 systems:

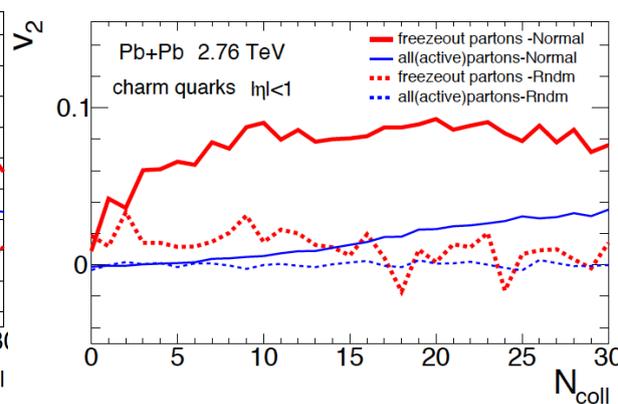
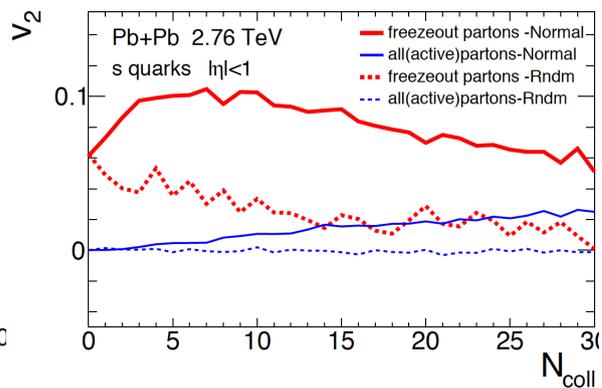
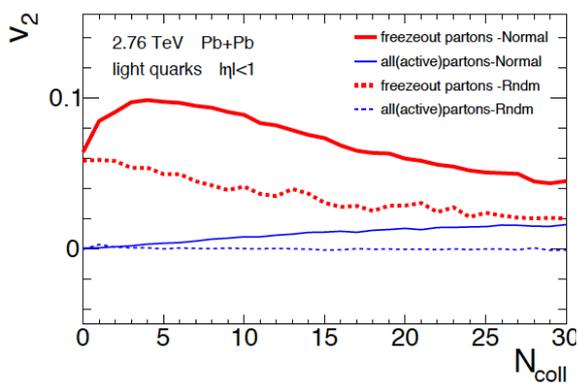
p+Pb  
5TeV  
b=0fm



Au+Au  
200GeV  
6.6-8.1 fm



Pb+Pb  
2.76TeV  
8fm



# The escape mechanism: *flavour dependence*

	<b>pPb</b> b=0fm	<b>AuAu</b> b=6.6-8.1fm	<b>PbPb</b> b=8fm
light	$\langle N_{\text{coll}} \rangle = 2.02$ $\langle v_2 \rangle_{\text{Rndm}} = 2.392\%$ $\langle v_2 \rangle_{\text{Norm}} = 3.279\%$ $\langle v_2 \rangle_{\text{Rndm}} / \langle v_2 \rangle_{\text{Norm}} = \mathbf{72.9\%}$	$\langle N_{\text{coll}} \rangle = 4.5$ $\langle v_2 \rangle_{\text{Rndm}} = 2.931\%$ $\langle v_2 \rangle_{\text{Norm}} = 4.468\%$ $\langle v_2 \rangle_{\text{Rndm}} / \langle v_2 \rangle_{\text{Norm}} = \mathbf{65.6\%}$	$\langle N_{\text{coll}} \rangle = 9.82$ $\langle v_2 \rangle_{\text{Rndm}} = 3.214\%$ $\langle v_2 \rangle_{\text{Norm}} = 7.562\%$ $\langle v_2 \rangle_{\text{Rndm}} / \langle v_2 \rangle_{\text{Norm}} = \mathbf{42.5\%}$
s-quark	$\langle N_{\text{coll}} \rangle = 2.54$ $\langle v_2 \rangle_{\text{Rndm}} = 1.894\%$ $\langle v_2 \rangle_{\text{Norm}} = 3.203\%$ $\langle v_2 \rangle_{\text{Rndm}} / \langle v_2 \rangle_{\text{Norm}} = \mathbf{59.1\%}$	$\langle N_{\text{coll}} \rangle = 5.45$ $\langle v_2 \rangle_{\text{Rndm}} = 2.266\%$ $\langle v_2 \rangle_{\text{Norm}} = 4.784\%$ $\langle v_2 \rangle_{\text{Rndm}} / \langle v_2 \rangle_{\text{Norm}} = \mathbf{47.4\%}$	$\langle N_{\text{coll}} \rangle = 11.14$ $\langle v_2 \rangle_{\text{Rndm}} = 2.23\%$ $\langle v_2 \rangle_{\text{Norm}} = 8.424\%$ $\langle v_2 \rangle_{\text{Rndm}} / \langle v_2 \rangle_{\text{Norm}} = \mathbf{26.5\%}$
c-quark	$\langle N_{\text{coll}} \rangle = 4.23$ $\langle v_2 \rangle_{\text{Rndm}} = 1.214\%$ $\langle v_2 \rangle_{\text{Norm}} = 2.139\%$ $\langle v_2 \rangle_{\text{Rndm}} / \langle v_2 \rangle_{\text{Norm}} = \mathbf{56.8\%}$	$\langle N_{\text{coll}} \rangle = 8.6$ $\langle v_2 \rangle_{\text{Rndm}} = 0.8455\%$ $\langle v_2 \rangle_{\text{Norm}} = 3.885\%$ $\langle v_2 \rangle_{\text{Rndm}} / \langle v_2 \rangle_{\text{Norm}} = \mathbf{22\%}$	$\langle N_{\text{coll}} \rangle = 15.48$ $\langle v_2 \rangle_{\text{Rndm}} = 0.6724\%$ $\langle v_2 \rangle_{\text{Norm}} = 7.923\%$ $\langle v_2 \rangle_{\text{Rndm}} / \langle v_2 \rangle_{\text{Norm}} = \mathbf{8.5\%}$

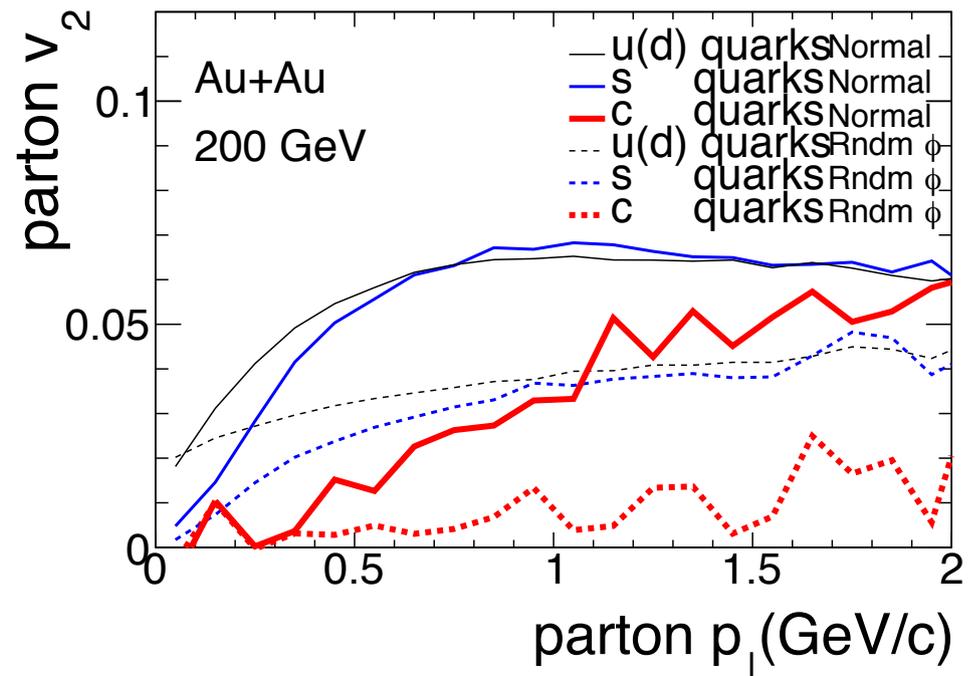
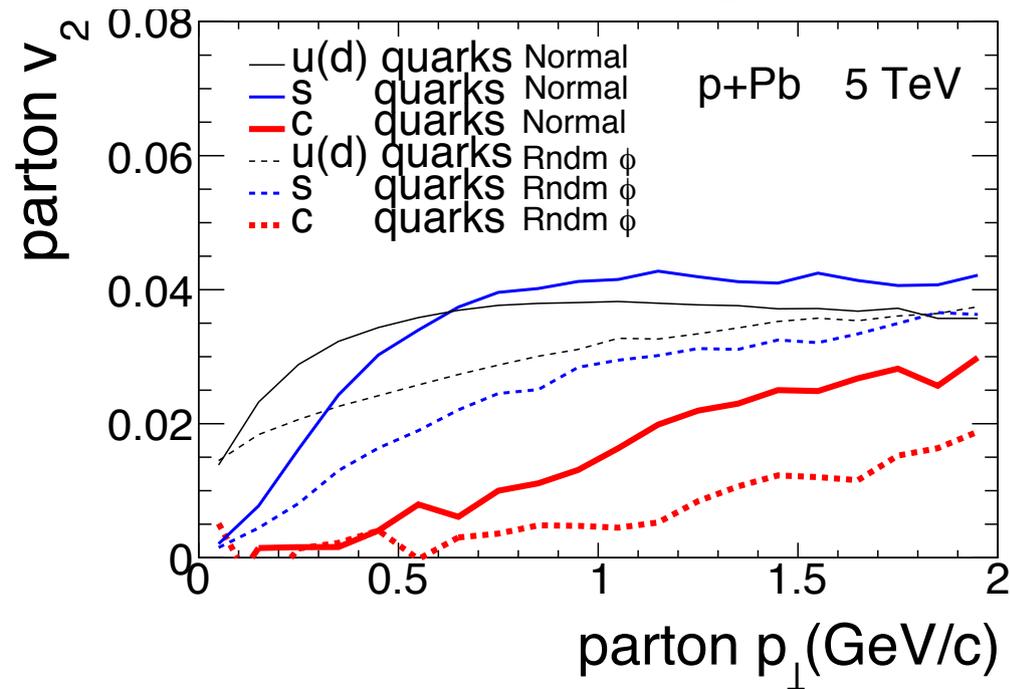
q

Q

System size/energy

Less from escape mechanism/  
more from collective flow

# The escape mechanism: *flavour dependence*



$v_2(P_T)$ :

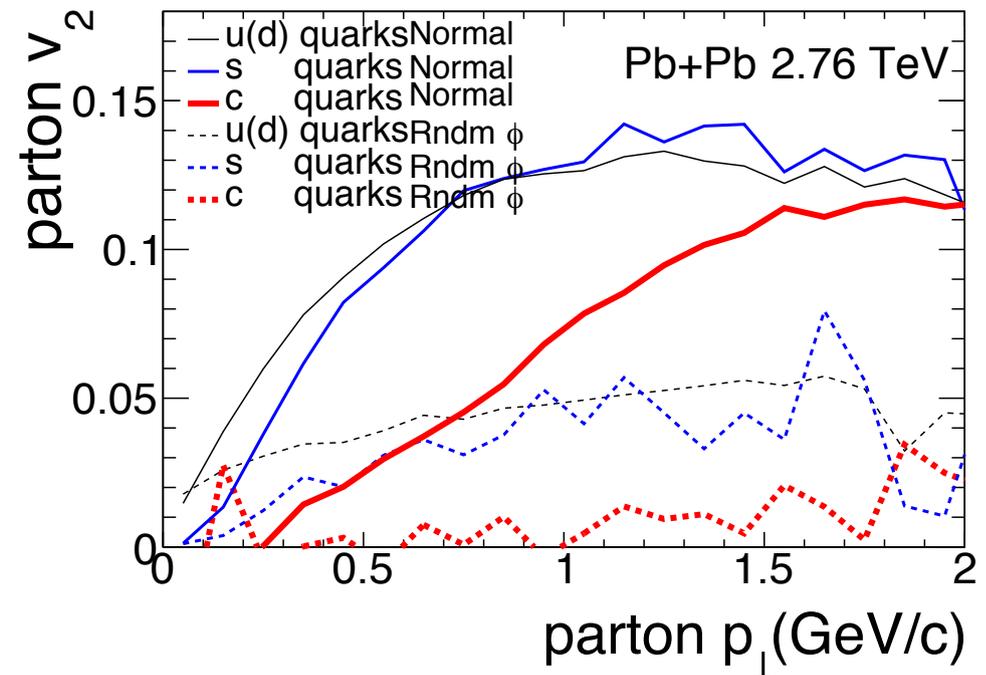
mass ordering at low  $P_T$ :

$$v_2^c < v_2^s < v_2^{ud}$$

this is partly responsible for  
the mass ordering of hadron  $v_2$  in  
Hanlin Li et al. PRC 93 (2016);  
arXiv:1604.07387

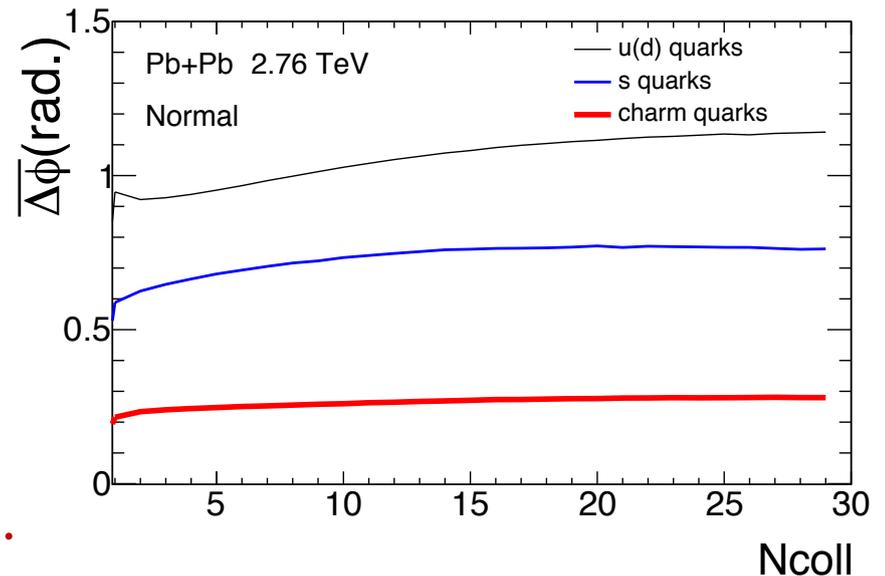
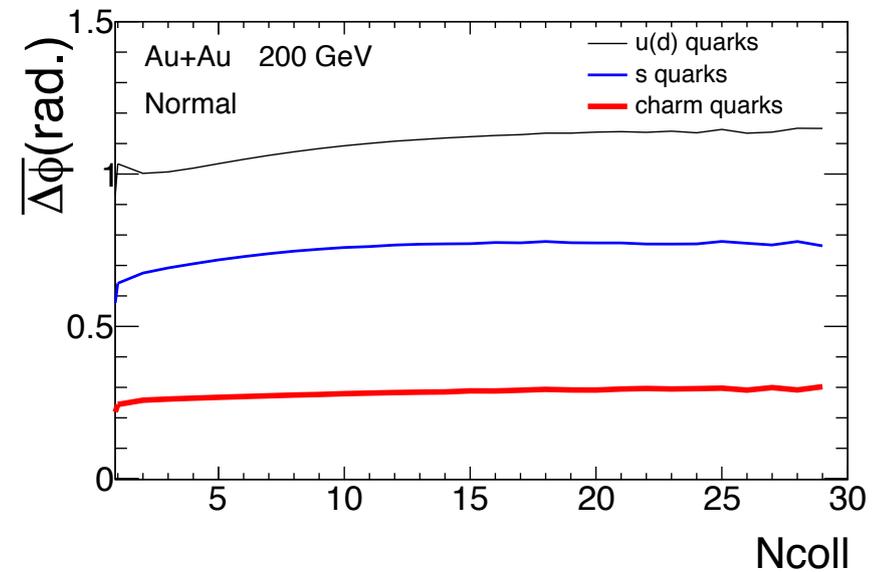
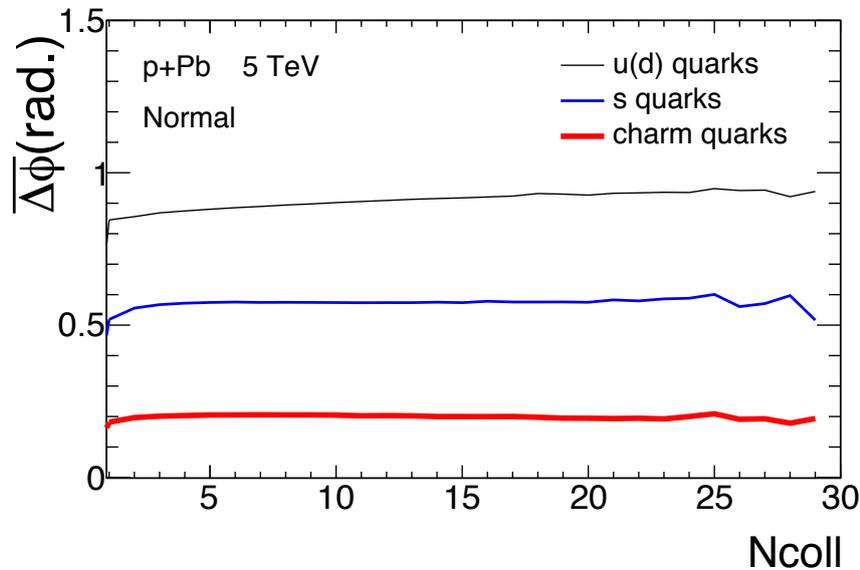
little mass effect at higher  $P_T$ :

$$v_2^c \sim v_2^s \sim v_2^{ud}$$



# The escape mechanism: *flavour dependence*

$\Delta\phi$ : change of azimuth due to one collision (the Ncoll-th collision):



Mass ordering on  
the mean parton deflection angle:

$$\overline{\Delta\phi}_c < \overline{\Delta\phi}_s < \overline{\Delta\phi}_{ud}$$

it is more difficult to deflect a heavier quark,  
so light quark flow  
& strong light-charm interactions  
are essential to generate significant charm  $v_2$ .

# Summary

We have followed the complete parton collision history to study  $v_2$  of light/strange/**charm** quarks in the AMPT model.

$\langle v_2 \rangle_{\text{random-}\phi} / \langle v_2 \rangle_{\text{normal}}$  ratio  $\sim$  fraction from pure escape:

	dAu@200GeV b=0 fm	pPb@5TeV b=0 fm	AuAu@200GeV b=6.6-8.1 fm	PbPb@2.76TeV b=8 fm
u/d	93%(all quarks)	72.9%	65.6%	42.5%
s		59.1%	47.4%	26.5%
<b>c</b>		<b>56.8%</b>	<b>21.8%</b>	<b>8.5%</b>

**$v_2$  of charm quarks** in large systems at high energies

mostly comes from collective flow (*not the escape mechanism*).

→ heavy quarks are more sensitive probes of collective flow & the medium.

Esha, Md. Nasim & Huang, JPG44 (2017); Greco's talk at QM2017.

serves as a comprehensive event generator for heavy ion collisions.

It aims to

evolve the system from initial condition to final observables;  
conserve energy/momentum/flavour/charge of each event,  
include particle productions of different flavours at different  $P_T$  &  $y$ ,  
keep non-equilibrium features and dynamics  
(e.g. intrinsic fluctuations and correlations).

It is also a test-bed of different ideas:

- Discovery of the triangular flow  $v_3$  Alver & Roland, PRC 81 (2010)
- Longitudinal (de)correlations of flows Pang et al. PRC 91 (2015), EPJA52 (2016)
- Flow may be dominated by anisotropic parton escape He et al. PLB753 (2016); ZWL et al. NPA 956 (2016)

Further efforts are needed to extend AMPT to heavy flavours, in order to simultaneously study light flavours, heavy flavours, including their interactions, to probe properties of the dense matter.