Conservation laws in net-particle fluctuation measurements: the effect on LQCD predictions of the QCD phase transition

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# The quark-hadron transition



1 μs after the Big Bang de-confined quarks and gluons transitioned into confined hadrons
 A key goal of HI collisions is to map out this transition

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#### Net-particle fluctuation measurements

# Recreating the Big Bang with heavy-ion collisions



Chemical freeze-out : all inelastic collisions between particle species cease
 Kinetic freeze-out : all elastic collisions between hadrons cease

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Net-particle fluctuation measurements

# Time evolution of heavy-ion collision



# Time evolution of heavy-ion collision



# Insight from theory



time

- There is the assumption that the system in HI collisions reaches thermalization
- Conserved quantities, Q<sub>i</sub> : net-charge, net-strangeness, and net-baryon number
- Described by GCE with partition function :

$$Z = \operatorname{Tr}\left[\exp\left(-\frac{H - \sum_{i} \mu_{i} Q_{i}}{T}\right)\right]$$
(1)

• Fluctuations of conserved quantities are described with susceptibilities :

$$\chi^{B,S,O}_{n_{B},n_{S},n_{O}} \equiv \frac{1}{VT^{3}} \frac{\partial^{n_{B}}}{\partial(\mu_{B}/T)^{n_{B}}} \frac{\partial^{n_{S}}}{\partial(\mu_{S}/T)^{n_{S}}} \frac{\partial^{n_{O}}}{\partial(\mu_{O}/T)^{n_{O}}} \ln Z$$
(2)

• Susceptibilities measure the response of the system to infinitesimal change in chemical potential  $\Rightarrow$  'Knobs' of the system to extract T and  $\mu_B$ 

# Theory in action

BNL-Bielefeld-CCNU



- Test of LQCD at  $\mu_B \approx 0$
- Smaller than in HRG for T > 150 MeV (right figure)
- Sellow boxes indicate the transition region,  $T_c \approx 154 \pm 9 \text{ MeV}$

#### What we know so far

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Can use experiment to map out the QCD phase transition!

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#### Fluctuations of conserved quantities in experiment



- Event-by-event particle multiplicity fluctuations are characterized by the cumulants of the event-by-event multiplicity distributions
- *N* is the net-particle number, i.e., particle *minus* antiparticle,  $\delta N = N \langle N \rangle$ , cumulants are :

$$C_{1} = \langle N \rangle$$

$$C_{2} = \langle (\delta N)^{2} \rangle$$
(3)
(3)
(4)

$$C_3 = \langle (\delta N)^3 \rangle \tag{5}$$

$$C_4 = \langle (\delta N)^4 \rangle - 3 \langle (\delta N)^2 \rangle^2 \tag{6}$$

• The skewness, S, and kurtosis,  $\kappa$ , can be extracted as follows :

$$S = \frac{C_3}{(C_2)^{\frac{3}{2}}}$$
(7)  

$$\kappa = \frac{C_4}{(C_2)^2}$$
(8)

# Generating fluctuations in experiment



- Fluctuations of net-baryons appear only inside finite acceptance
- The fluctuations can be considered as meaningfully observables if the following is true (GCE) :

$$\Delta Y_{accept} \gg \Delta Y_{corr} \tag{9}$$

$$\Delta Y_{total} \gg \Delta Y_{accept} \gg \Delta Y_{kick} \tag{10}$$

- $\Delta Y_{accept}$  : acceptance of measurement
- $\Delta Y_{total}$  : full phase space
- $\Delta Y_{corr}$  : rapidity correlation length
- $\Delta Y_{kick}$  : rapidity shift due to hadronization

#### Connecting theory to experiment

• The cumulants of the conserved charges are related to the susceptibilities :

$$C_{ijk}^{BOS} = VT^3 \chi_{ijk}^{BOS}(T, \mu_B, \mu_Q, \mu_S)$$
<sup>(11)</sup>

• Since the volume is unknown in experiment, take ratios, e.g. :

$$\frac{C_2}{C_1} = \frac{\sigma^2}{M} = \frac{\chi_2}{\chi_1}$$
(12)
$$\frac{C_3}{C_2} = S\sigma = \frac{\chi_3}{\chi_2}$$
(13)
$$\frac{C_4}{C_2} = \kappa\sigma^2 = \frac{\chi_4}{\chi_2}$$
(14)
$$\frac{C_4}{C_3} = \frac{\kappa\sigma}{S} = \frac{\chi_4}{\chi_3}$$
(15)

• Focus of this talk will be on the highlighted ratio

## Theory and experiment in action

▶ S. Borsanyi et al ▶ STAR collab.



#### Limitations with connecting theory to experiment

- Some assumptions and approximations : The net-particle multiplicity distribution from experiment follows a Poisson PD
  - $\implies$  Skellam baseline :  $C_n = \langle N_B \rangle + (-1)^n \langle N_{\bar{B}} \rangle$
  - $\implies$  Proxies : e.g. net-proton  $\approx$  net-baryon
- Global charge conservation : experiment mimics GCE by analyzing a subset of the particles in the final state
  - $\implies$  acceptance cuts can lead to global conservation laws

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# Net-proton fluctuation measurement from ALICE



ALICE collab.



ALICE collab.





## Baryon number conservation in fluctuation measurements

P. Braun-Munzinger et al ALICE collab.



The data is best described with the assumption of global baryon number conservation  $\checkmark$  Consistent with LQCD prediction of a Skellam behavior for  $\kappa_2$  of net-baryons after accounting for baryon number conservation

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# **Net-** $\Lambda$ fluctuation measurement from STAR

# Beam energy dependence of net-Λ fluctuations

STAR collab. R. Bellwied et al





#### Beam energy dependence of net-A fluctuations

STAR collab. R. Bellwied et al





A carries strangeness : test for presence of quark flavor dependent freeze out

Need higher order cumulants to distinguish FO

#### Rapidity dependence of net-A cumulants

STAR collab.



- Acceptance factor  $\alpha$  includes strangeness
- $\blacksquare \approx$  works for 19.6 GeV data
- Coupling of strangeness and baryon number conservation is unclear

Net-particle fluctuation measurements

- Susceptibilities from lattice QCD calculations are directed linked to net-particle fluctuation measurements from experiment
  - $\implies$  Net-proton fluctuation measurement of the second cumulant from ALICE is consistent with LQCD predictions after accounting for baryon number conservation
  - $\implies$  Coupling between strangeness and baryon number requires more studies

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  - $\implies$  Coupling between strangeness and baryon number requires more studies
- Higher order cumulants of net-proton at LHC and of net-A at RHIC will take advantage of higher statistics in upcoming runs



Third order cumulant dependence on acceptance in simulation P. Braun-Munzinger, A. Rustamov, J. Stachel